

Formation of columnar baroclinic vortices in thermally stratified nonlinear spin-up*

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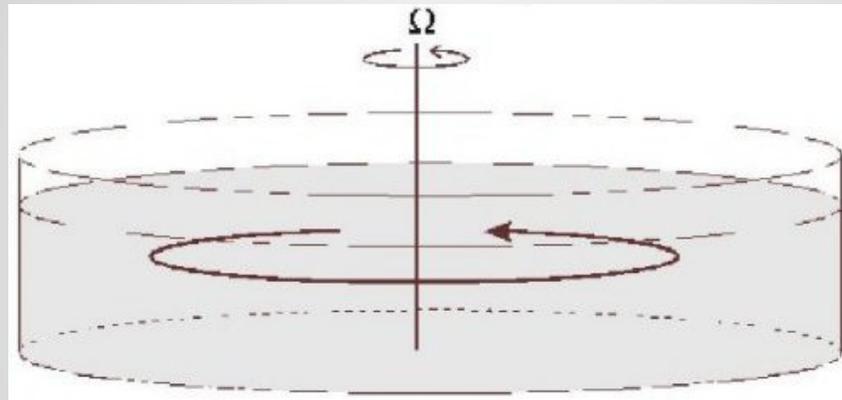
Roberto Verzicco at Universita' di Roma 'Tor Vergatta', Italy.

Outline

- Overview
 - What are spin-up and baroclinic flows?
- Background
 - Laboratory experiments
 - Numerical simulations
- Results
- Conclusions
- Problems to consider in the future

Overview

Definition of spin-up/spin-down: an impulsive change of the rotation rate of a rigid container filled with a liquid



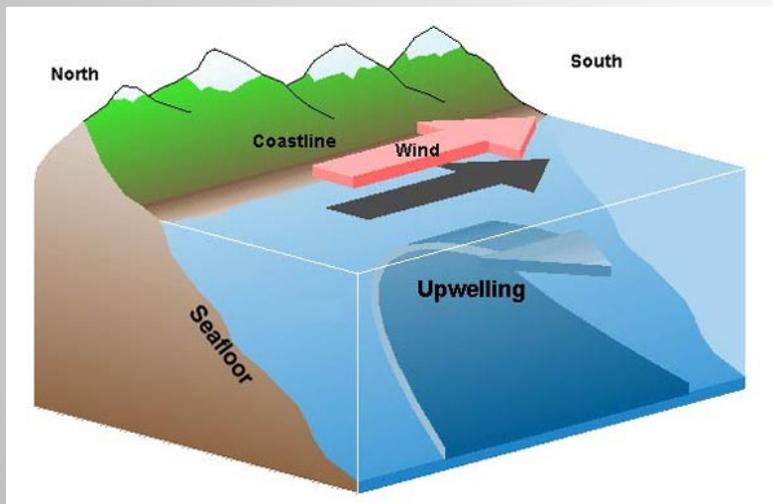
- fluid assumed to be in solid body rotation
- rotation rate of the container modified: increased (**spin-up**) or decreased (**spin-down**)
- adjustment of the primary circulation to the new rotation rate

Background

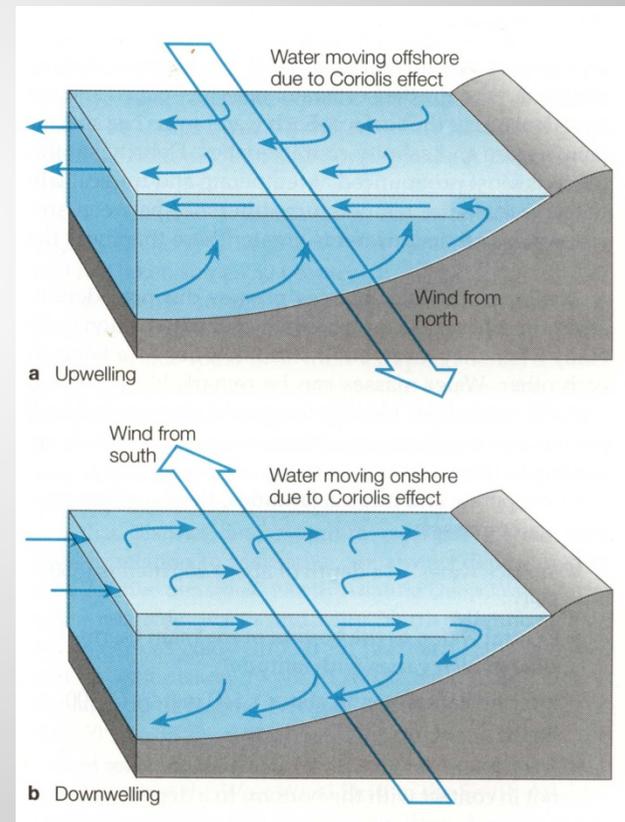
- Existence of separated time scales: Ekman number $E = \nu / \Omega H^2$
 - Ekman boundary layer development: $\tau_e = \Omega^{-1}$
 - Spin-up due to secondary circulation: $\tau_s = \Omega^{-1} E^{-1/2}$
 - Viscous dissipation of residual motions: $\tau_v = \Omega^{-1} E^{-1}$
- Large body of work on the axisymmetric stratified spin-up, but not much on the non-axisymmetric counterpart
 - Numerical models: axisymmetric in nature
 - Experimental investigations: mostly in small facilities

Geophysical flows

Response mechanism to external forcing which arises in oceans

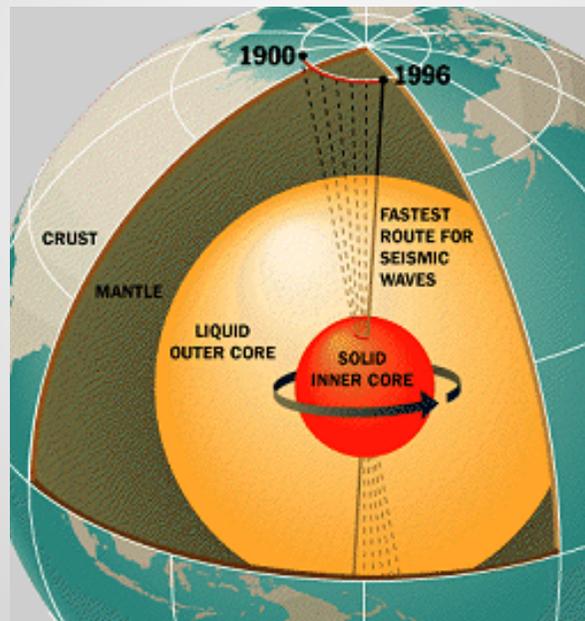


Ekman transport moves surface waters away from the coast; these are replaced by denser water that wells up from below



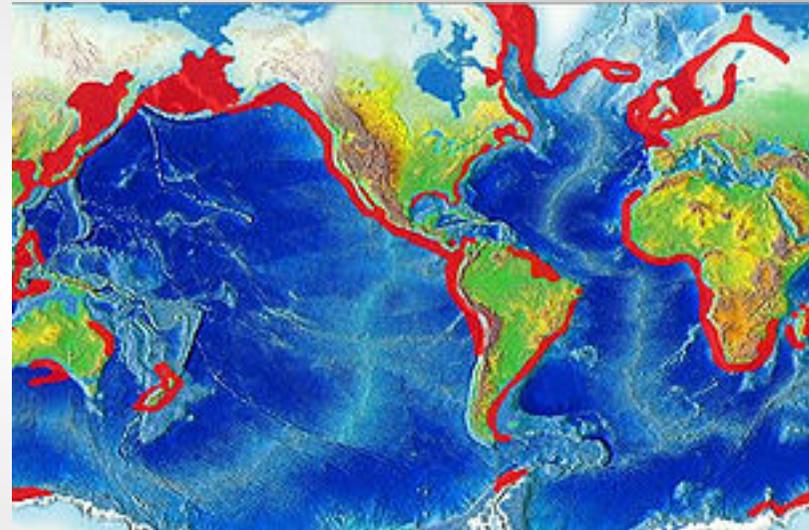
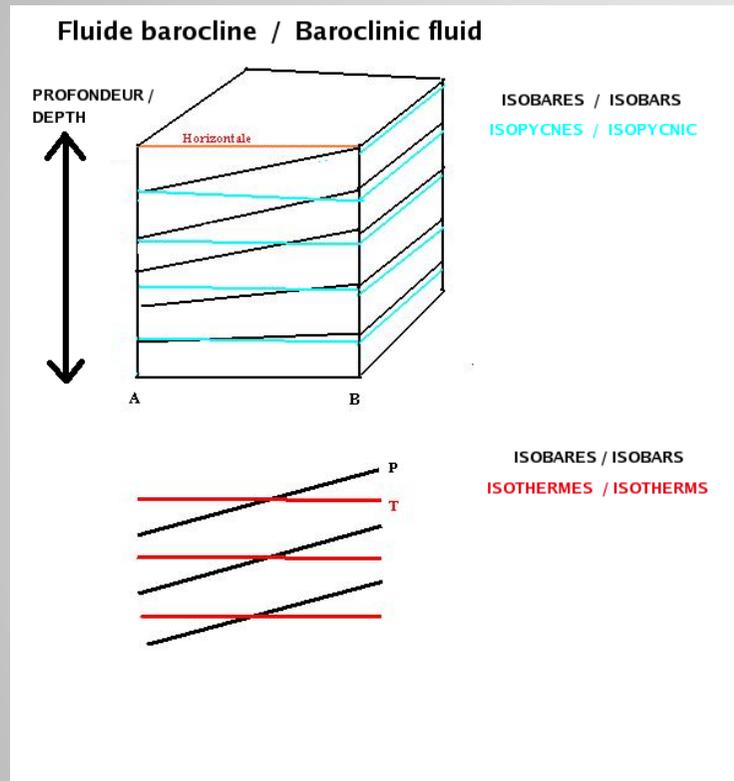
Magnetohydrodynamic flows

Response mechanism to magnetic forcing which arises in the inner core of the earth



Example: Mercury is known to have an at least partially fluid core, and during its orbit its rotation rate increases and decreases periodically.

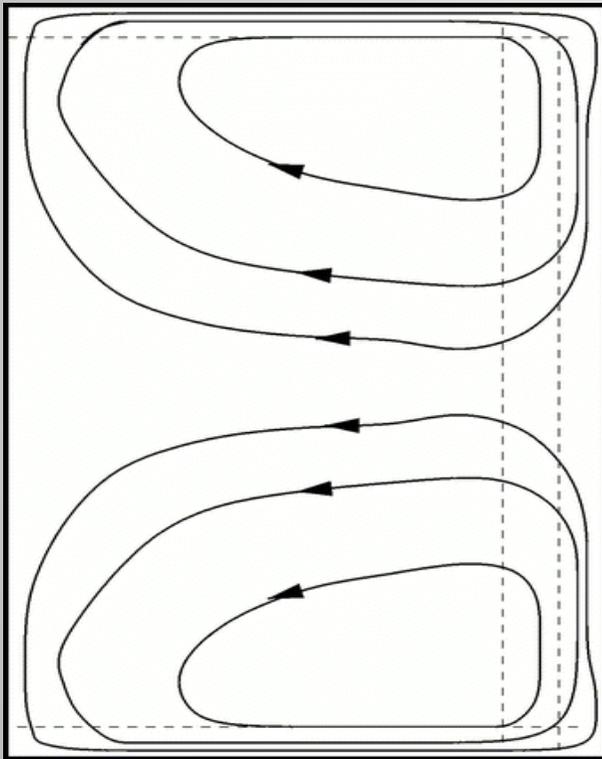
Baroclinic fluid



Spin-up homogeneous fluid

centerline

sidewall



Dashed lines – Ekman layers & sidewall Stewartson layers (all of which transport mass)

Stewartson (1957)

Greenspan & Howard (1963)

$$\varepsilon = \frac{\Delta\Omega}{\Omega} \quad E = \frac{\nu}{\Omega H^2}$$

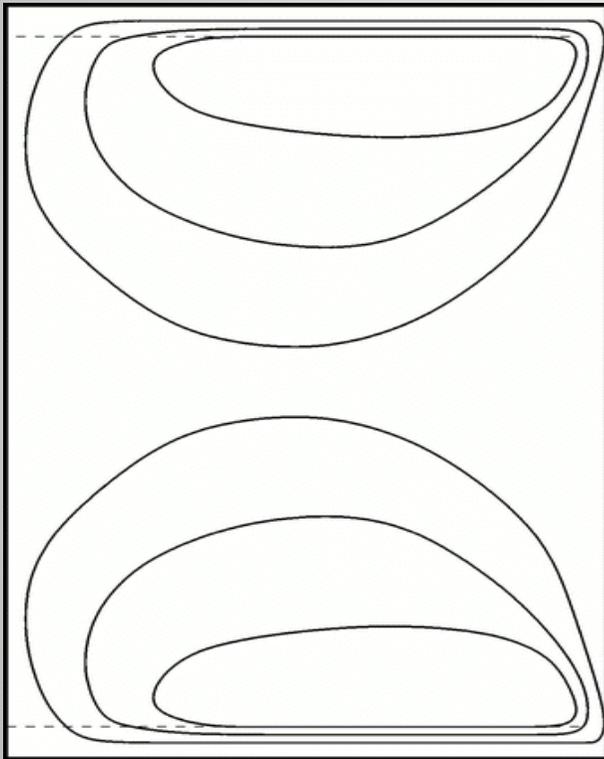
$$\tau_s = \Omega^{-1} E^{-1/2}$$

Streamlines for homogenous
spin-up

Spin-up stratified fluid

centerline

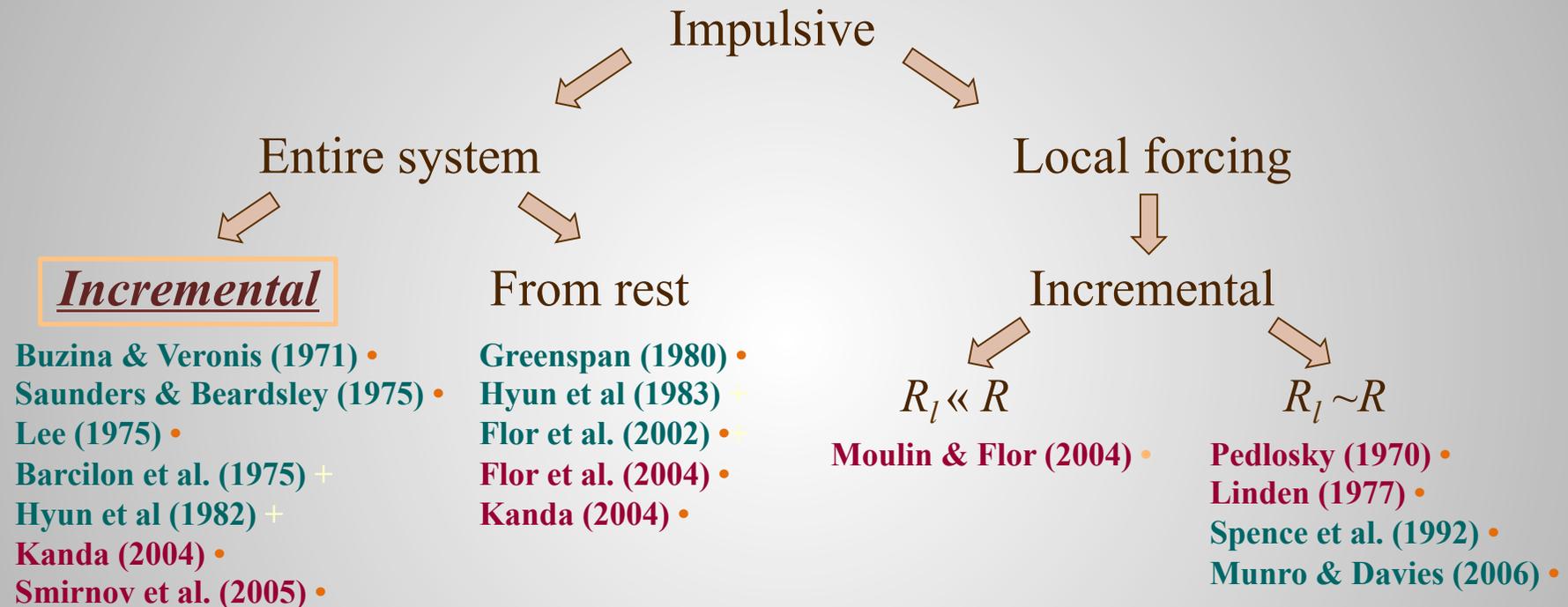
sidewall



Streamlines for stratified
spin-up

- Inhibition of vertical motions by buoyancy force
- Presence of horizontal density gradients
- Baroclinic instability ($Bu < 1$)
 $Bu \equiv N^2 H^2 / \Omega^2 R^2$
- Vertical shear of the horizontal velocity
- Ekman layers have the same structure as in homogeneous case

Stratified spin-up



+ Smirnov, Pacheco and Verzicco *Physics of Fluids* (2010) +

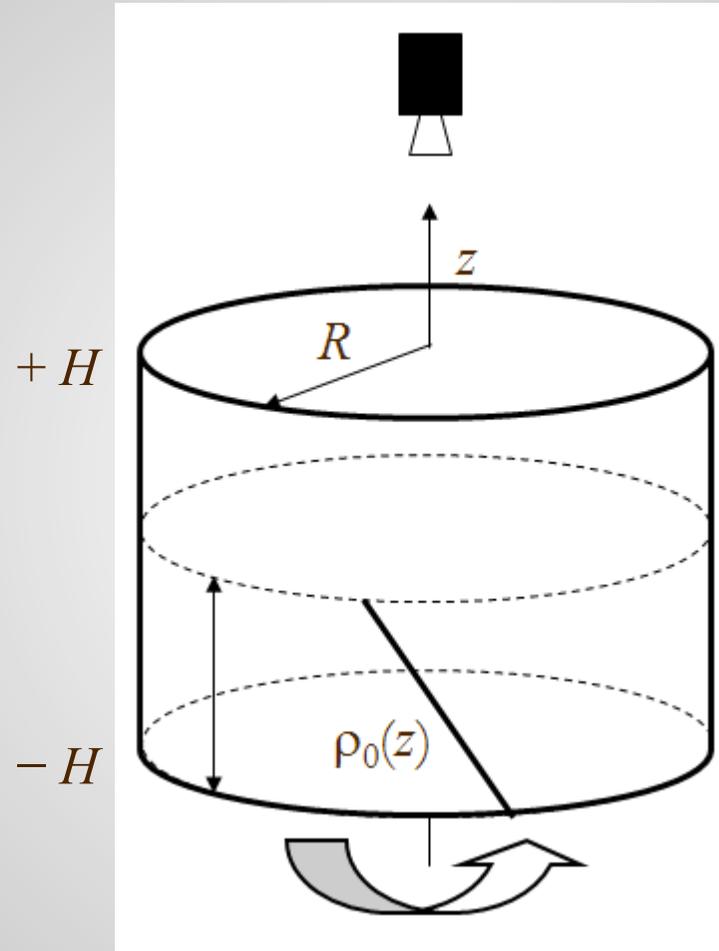
axisymmetric flows
non-axisymmetric flows

• experimental study
+ numerical work

R_l = radius of lid
 R = radius of horizontal boundaries

Experiments of Smirnov et al. (2005)

Salinity stratification



$$\Omega(1 - \varepsilon) \rightarrow \Omega \quad \varepsilon = \Delta\Omega/\Omega$$

Stratified spin-up

Parameters:

$$\varepsilon = 0.24 \quad B_u = 0.24$$

$$f = 0.5 \text{ s}^{-1} \quad R_d = 22 \text{ cm}$$

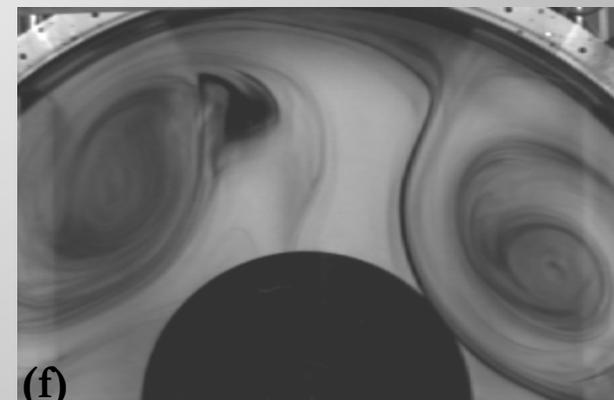
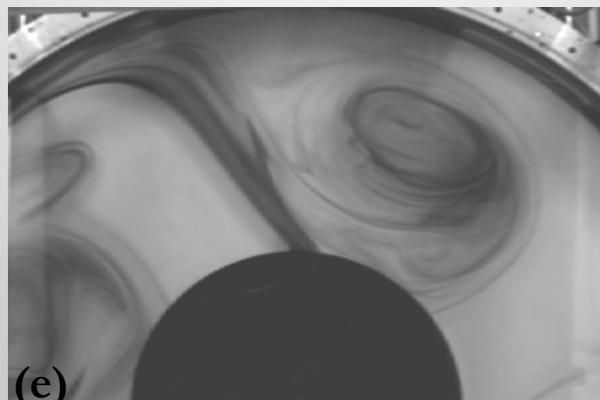
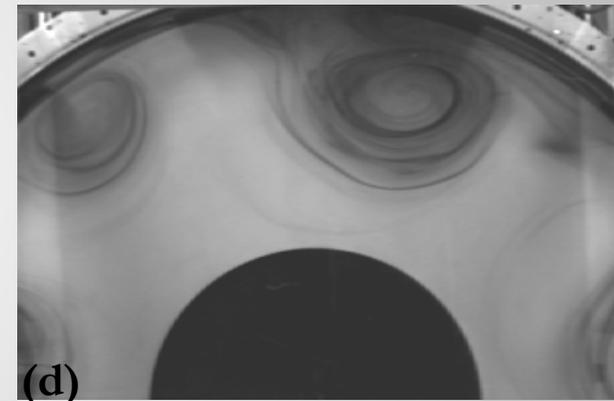
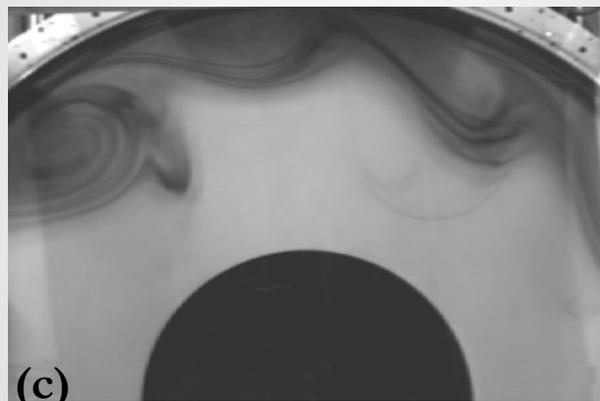
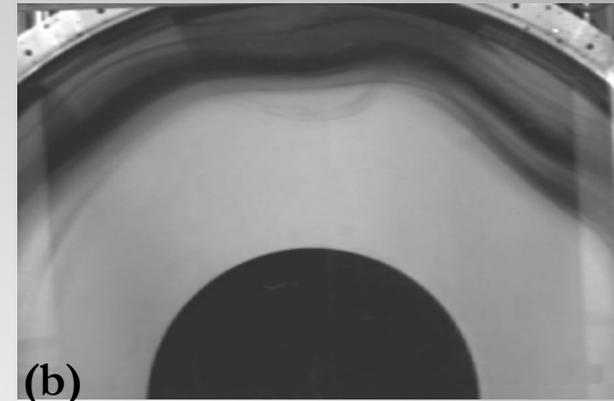
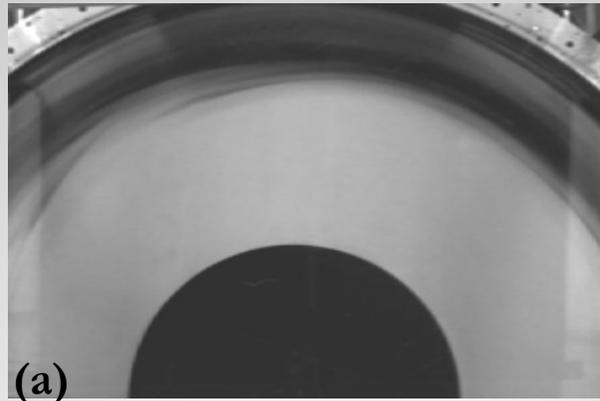
$$\Omega^{-1}E^{-1} \sim 4.5 \text{ hours}$$

t/T

$$(a) 11.6 \quad (b) 20.2$$

$$(c) 30 \quad (d) 32.6$$

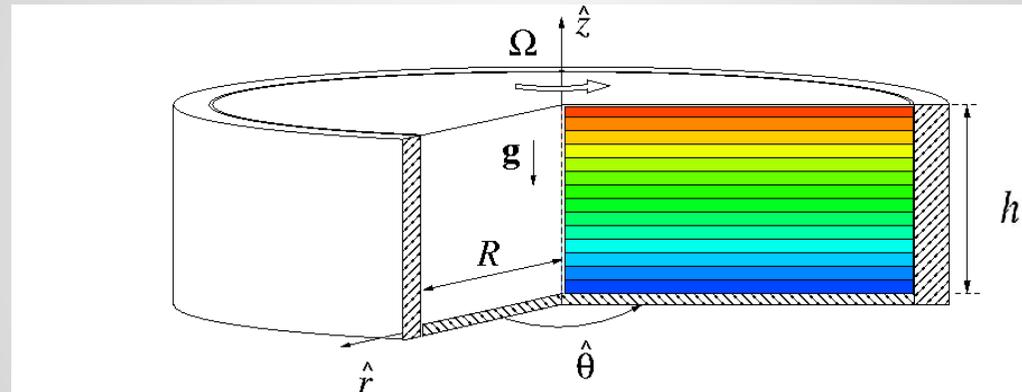
$$(e) 39.2 \quad (f) 43.2$$



Description of the problem

Consider the flow in a cylindrical region with radius R and height h .

At $t = 0$ the system is accelerated from its initial state of solid body rotation with angular velocity Ω_i to a new rotation rate $\Omega = \Omega_i + \Delta\Omega$.



Governing equations

$$\frac{D\vec{u}}{Dt} + 2\vec{\Omega} \times \vec{u} = -\frac{1}{\rho_o} \nabla p - \frac{\rho}{\rho_o} g \vec{e}_z + \frac{\rho}{\rho_o} \vec{\Omega}^2 r \vec{e}_r + \nu \nabla^2 \vec{u},$$

$$\frac{DT}{Dt} = \kappa \nabla^2 T, \quad \nabla \cdot \vec{u} = 0, \quad \rho = \rho_o (1 - \alpha T).$$

Initial conditions

$$u = w = 0, v = -\Delta\Omega r$$

$$\frac{\partial T}{\partial z} = \frac{\Delta T}{h}$$

First 3D numerical simulations of spin-up with thermal stratification and adiabatic/isothermal endwalls*

Boundary conditions (BC):

$$u = v = w = \frac{\partial T}{\partial r} = 0, \quad \text{at } r = R$$

$$u = v = w = 0, \quad \text{at } z = 0$$

$$\frac{\partial u}{\partial z} = \frac{\partial v}{\partial z} = w = 0, \quad \text{at } z = h$$

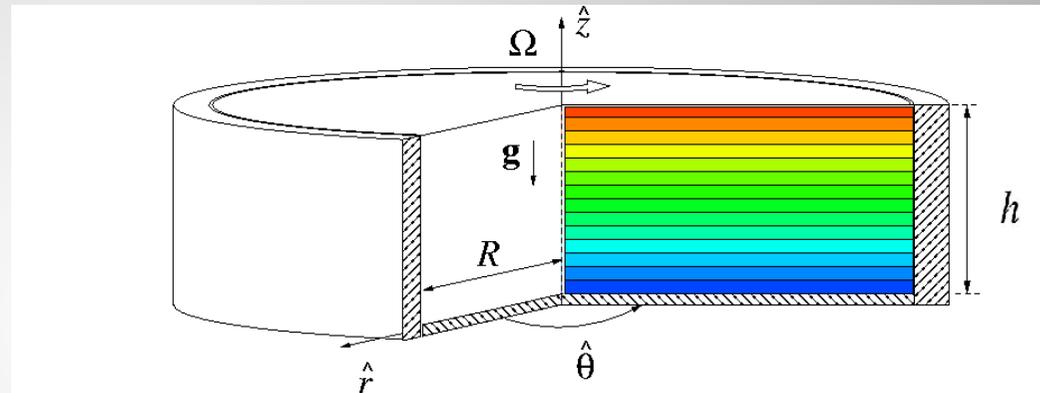
Adiabatic (Neumann) BC

$$\frac{\partial T}{\partial z} = 0, \quad \text{at } z = 0, h$$

Isothermal (Dirichlet) BC

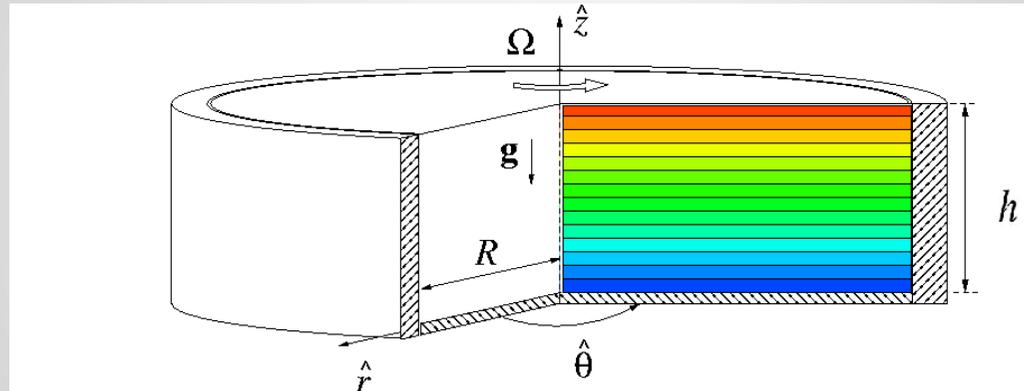
$$T = T_b, \quad \text{at } z = 0,$$

$$T = T_t, \quad \text{at } z = h$$



Navier Stokes equations

$$\begin{aligned}(\partial_t + \mathbf{u} \cdot \nabla) \mathbf{u} &= \nabla p + B^2 \Theta \mathbf{e}_z + 2\mathbf{u} \times \mathbf{e}_z - F B^2 \Theta r \mathbf{e}_r + E \nabla^2 \mathbf{u}, \\ (\partial_t + \mathbf{u} \cdot \nabla) \Theta &= \sigma^{-1} E \nabla^2 \Theta, \quad \nabla \cdot \mathbf{u} = 0,\end{aligned}$$



Ekman number: $E = \nu / \Omega h^2$,

Froude number: $F = \Omega^2 h / g$,

Burger number: $B = N / \Omega$,

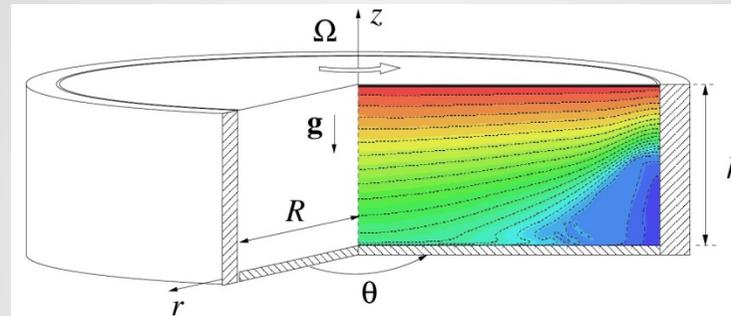
Prandtl number: $\sigma = \nu / \kappa$,

aspect ratio: $\Gamma = R / h$,

Rossby number: $\epsilon = \Delta \Omega / \Omega$,

Results of numerical simulations

Look at the isotherms on the planes $\theta = 0 - \pi$

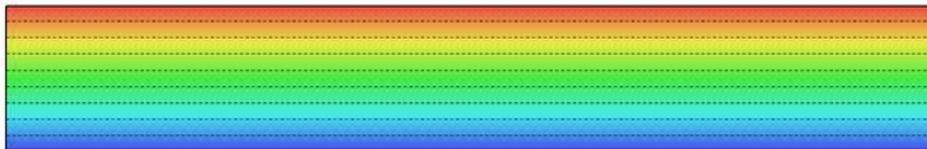


$$R/h=3.3, \varepsilon=0.73$$

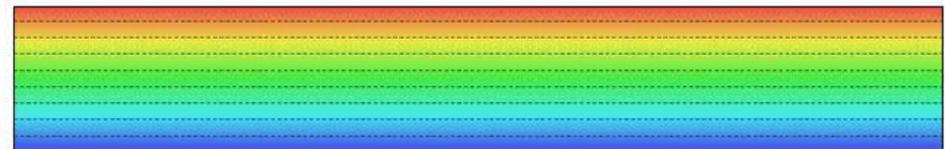
Adiabatic BC

Isothermal BC

Number of rotation periods = 1



Neumann boundary conditions for temperature



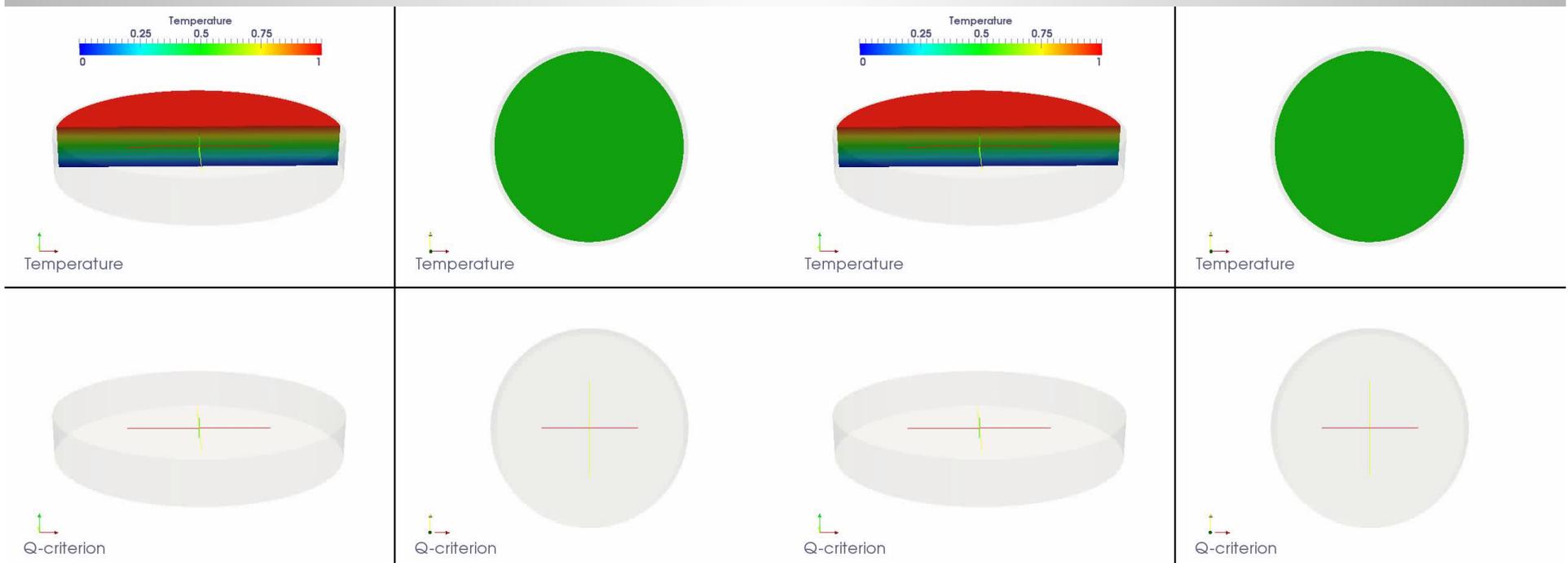
Dirichlet boundary conditions for temperature

Vortex structure

Vortex structure identified by the iso-surfaces of $Q = 0$ (Hunt et al. 1988)
colored by temperature

Adiabatic BC

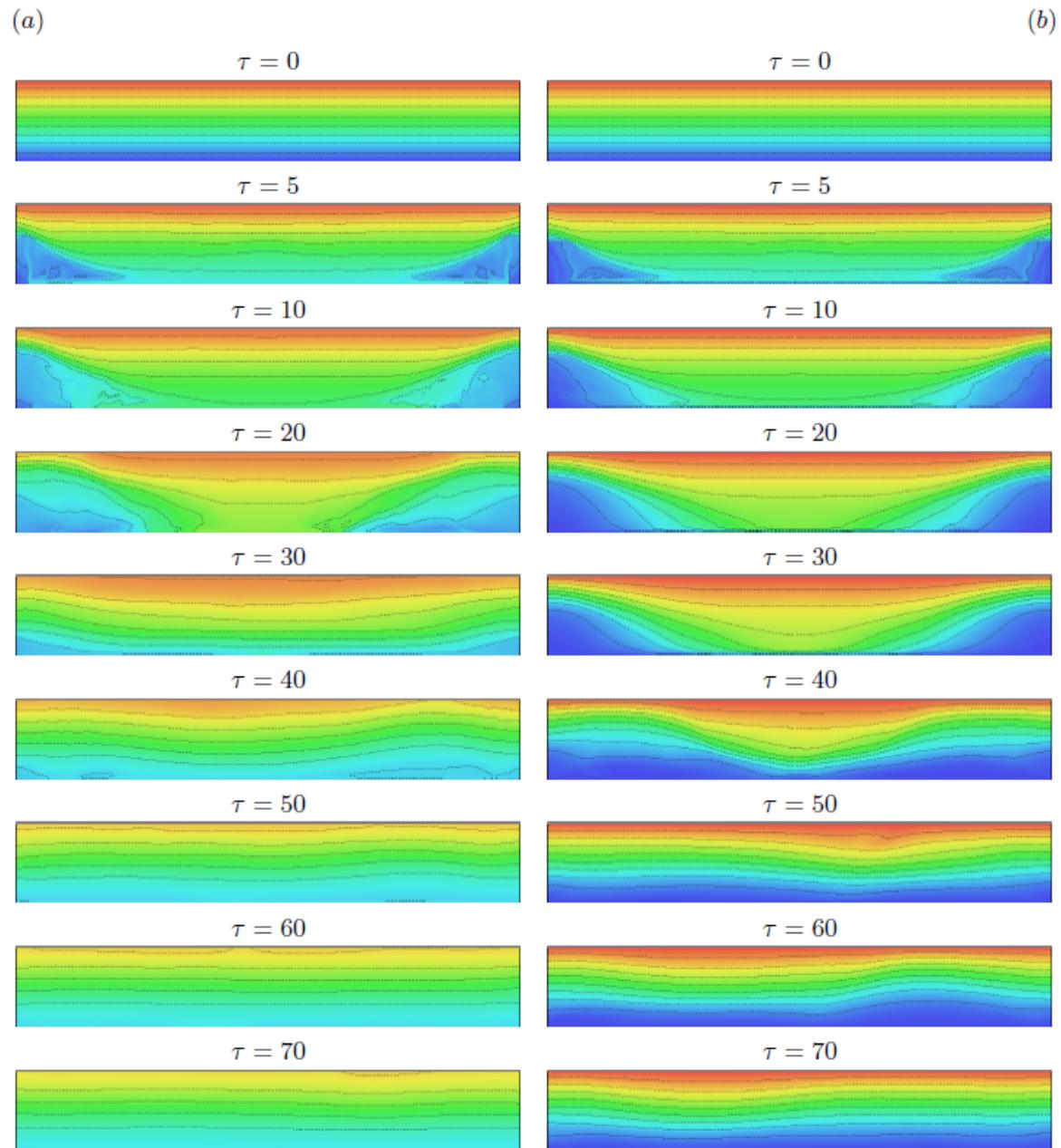
Isothermal BC



Adiabatic BC

$R/h=3.3, \epsilon=1$

Isothermal BC



Observations:

- Columnar eddies seen in experiments are reproduced in a numerical setting
- The emergence of these columnar eddies are influenced by the type of boundary condition (adiabatic/isothermal)
- Forced stratification suppresses the instability, i.e. isothermal BC are more stable than adiabatic BC
- Columnar vortices with isothermal boundary conditions appear at a later time compared to those with adiabatic boundary conditions

Why?

Hypothesis: production of baroclinic vorticity is responsible for the formation of vertical structures

Vorticity: $(\partial_t + \mathbf{u} \cdot \nabla)\boldsymbol{\omega} = (\boldsymbol{\omega} + 2\mathbf{e}_z) \cdot \nabla \mathbf{u} - F B^2 \nabla(r\Theta) \times \mathbf{e}_r + B^2 \nabla\Theta \times \mathbf{e}_z + E \nabla^2 \boldsymbol{\omega}$

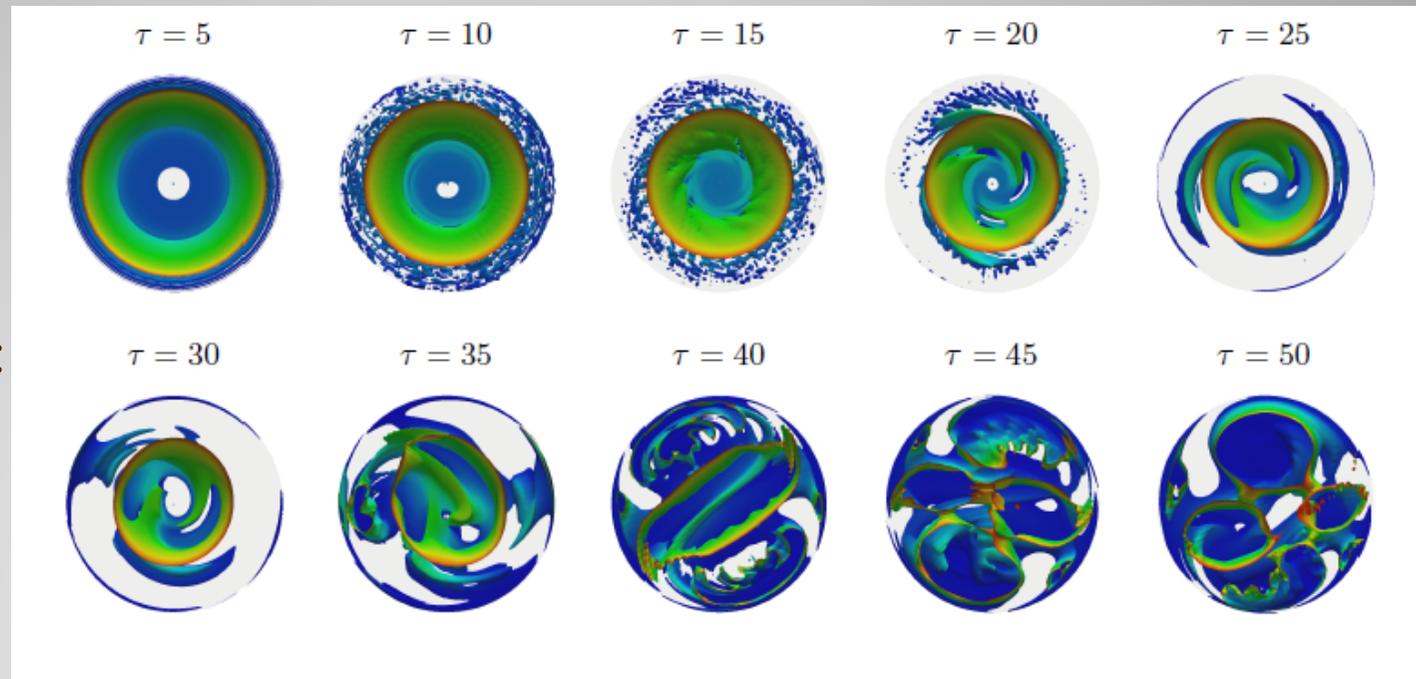
Baroclinic vorticity: $\boldsymbol{\omega}_b = \nabla\Theta \times \mathbf{e}_z = \left(\frac{1}{r} \frac{\partial}{\partial \theta} \mathbf{e}_r - \frac{\partial}{\partial r} \mathbf{e}_\theta \right) \Theta.$

Radial variation cannot produce vertical structures!

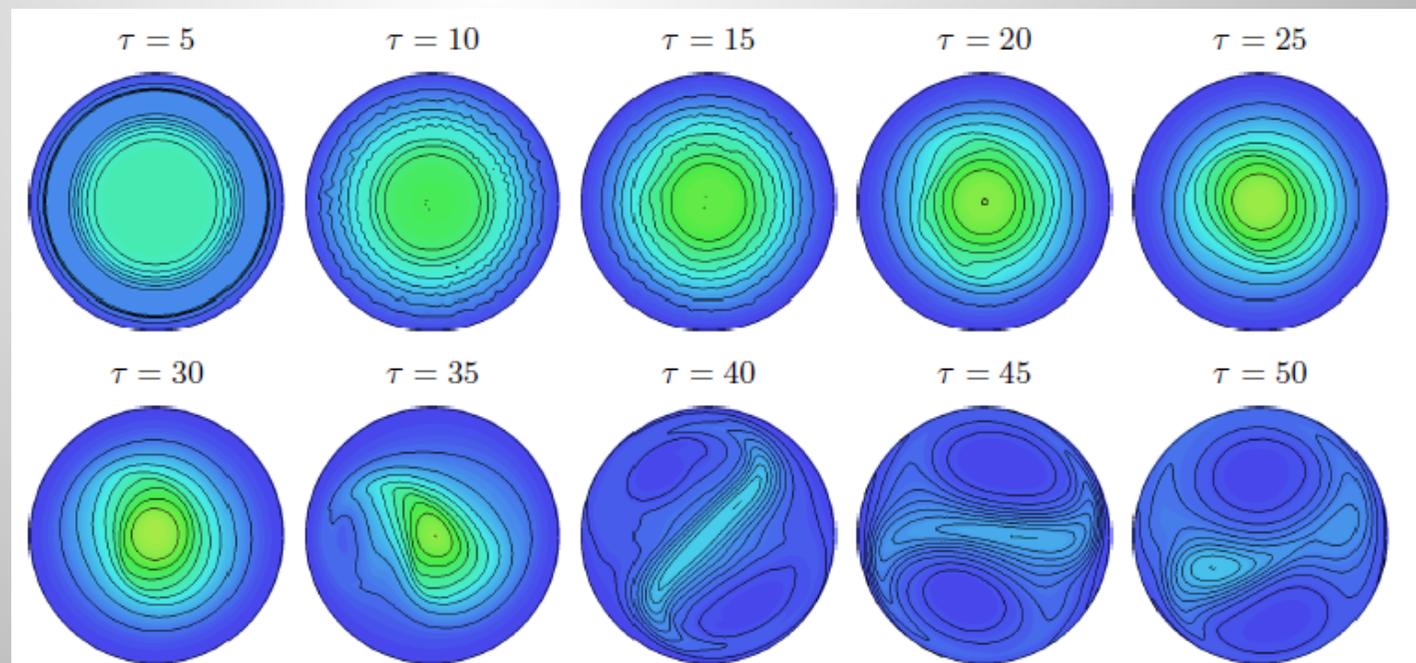
Azimuthal variation cannot produce vertical structures!

Feedback mechanism

Vortex structure:



Contours of temperature ($z=0.15$):



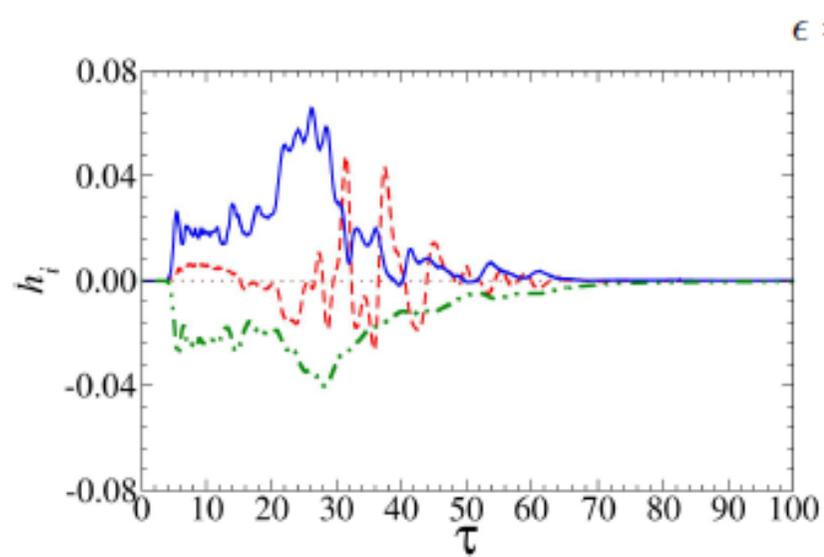
Kinetic energy growth rate of the azimuthal disturbance:

$$\frac{de}{dt} = \frac{d}{dt} \int_V \frac{1}{2} |\vec{u}'|^2 dV = \underbrace{-\int_V \vec{u}' \cdot (\vec{u}' \cdot \nabla' \vec{u}) dV}_{h_1} + \underbrace{B^2 \int_V \Theta u'_z dV}_{h_2} - \underbrace{FB^2 \int_V \Theta r u'_r dV}_{h_3} - \underbrace{E \int_V |\nabla' \vec{u}|^2 dV}_{h_4}$$

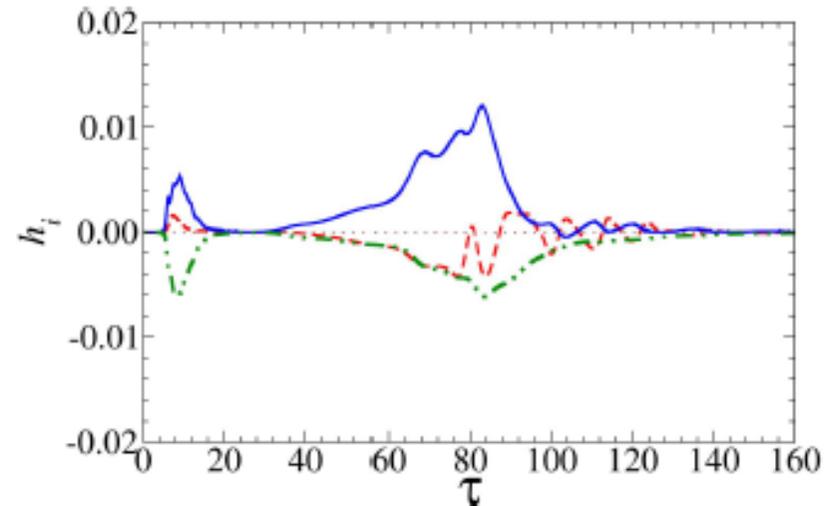
- h_1 : shear of the mean axisymmetric flow (barotropic production)
- h_2 : conversion of gravitational potential energy (baroclinic production)
- h_3 : conversion of centrifugal potential energy
- h_4 : viscous dissipation

Kinetic energy growth rate of the azimuthal disturbance

Adiabatic BC



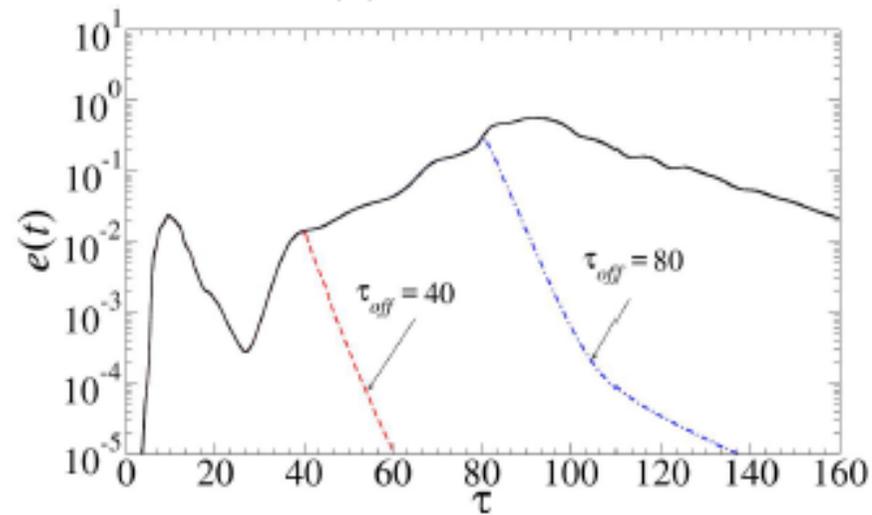
Isothermal BC



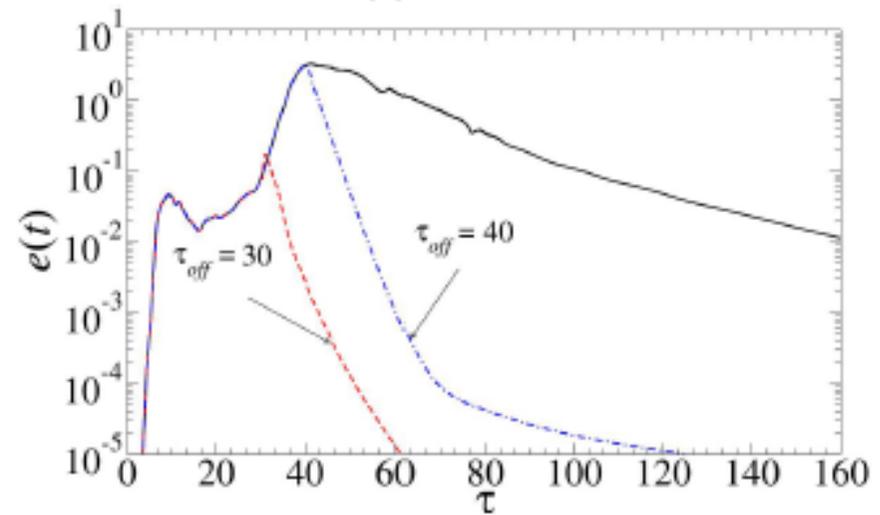
- h_1 : shear of the mean axisymmetric flow (barotropic production)
- h_2 : conversion of gravitational potential energy (baroclinic production)
- h_3 : conversion of centrifugal potential energy
- h_4 : viscous dissipation

Turn off the baroclinic term

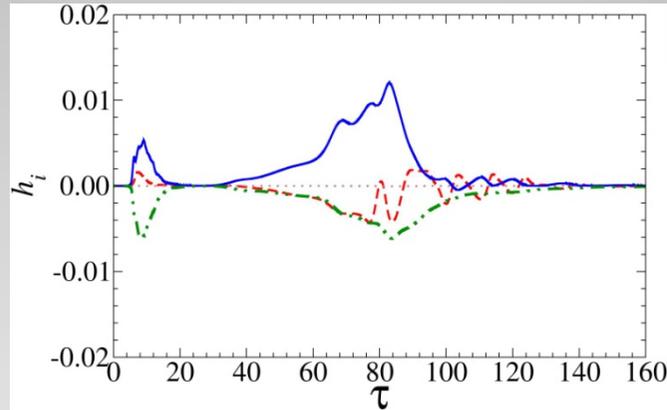
(a) $\epsilon = 0.73$



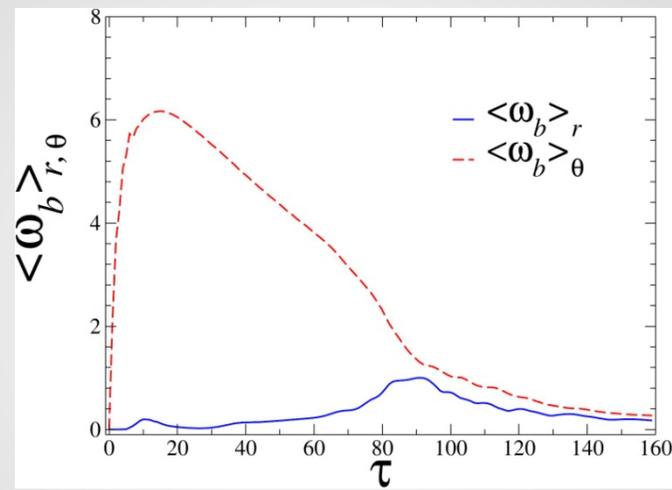
(b) $\epsilon = 1$



- h_i terms of the energy equation of azimuthal disturbances

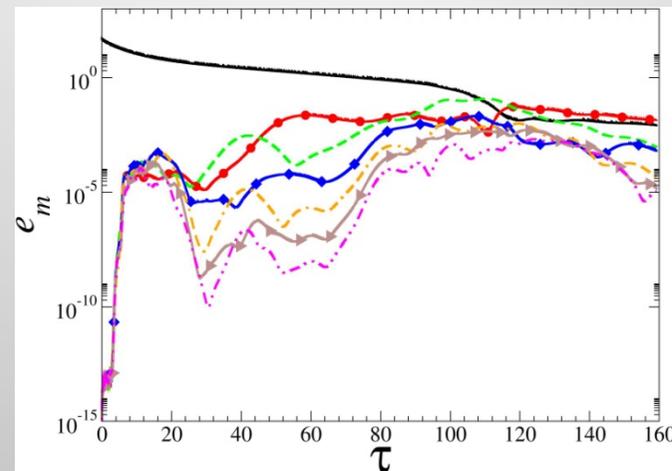


- Total baroclinic vorticity components



$$\omega_b = \nabla \Theta \times e_z = \left(\frac{1}{r} \frac{\partial}{\partial \theta} e_r - \frac{\partial}{\partial r} e_\theta \right) \Theta.$$

- Modal energy e_m



Conclusions*

- Radial variations of baroclinic vorticity produce an unstable system
- The azimuthal variations of baroclinic vorticity responsible for the formation of columnar vortices at late stages of flow development via feedback mechanism (enhances axial vorticity and increases vertical shear)
- Shear-free lateral boundary is more unstable
- Elliptical instability not important in this geometry

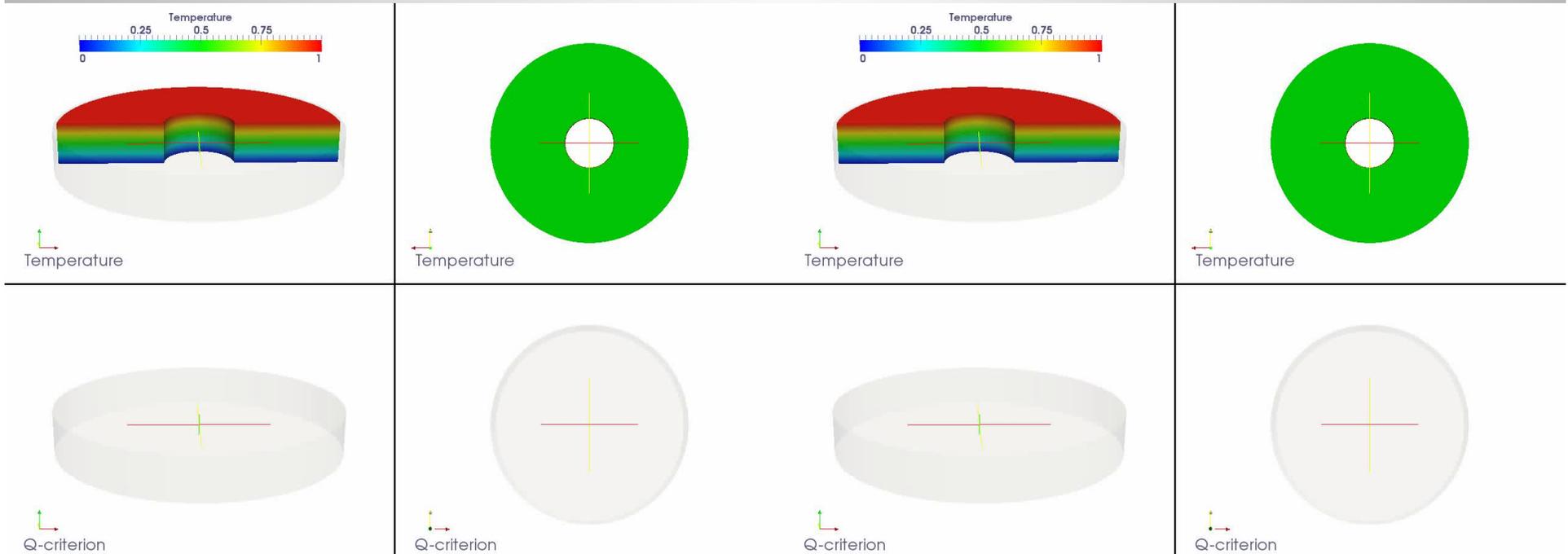
What other aspects of spin-up/down can be studied?

- Study in parameter space ($\Gamma, B, \varepsilon, \sigma$)
- Spin-up with sloping bottom (Ekman arrest)
- Uneven bottom boundary (wavy bottom)
- Heating/cooling on horizontal walls
- Non-Newtonian flows (lava flows)
- Different geometries: annulus, cone, half cone, half cylinder, etc.
- Spin-down
- Libration
- Turbulence!!!

Short annulus

Neumann BC

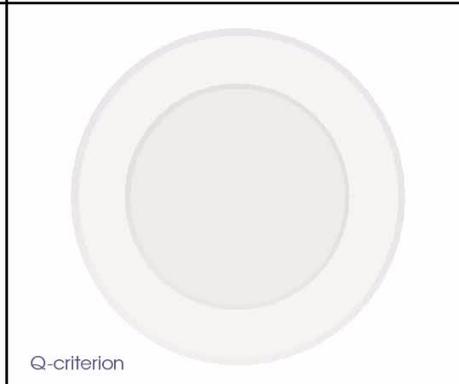
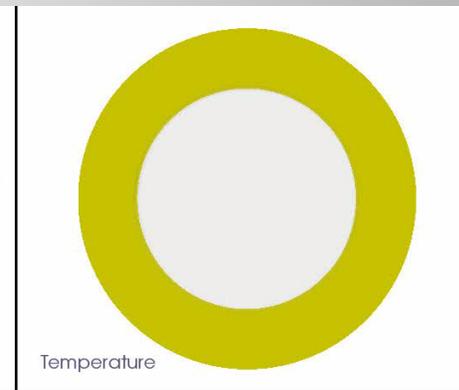
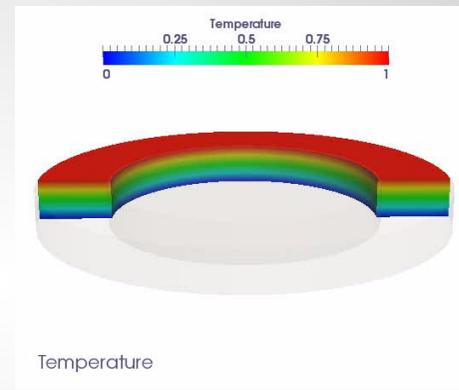
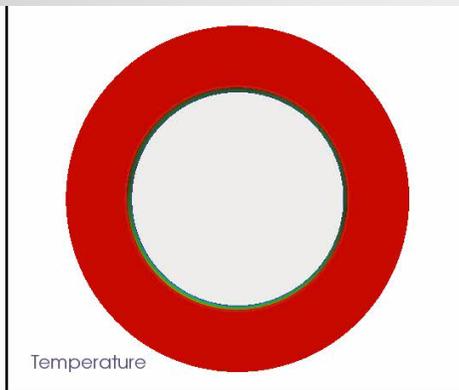
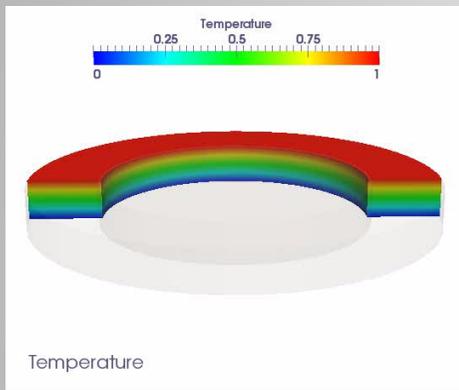
Dirichlet BC



Thank you!

Neumann BC

Dirichlet BC



Ongoing

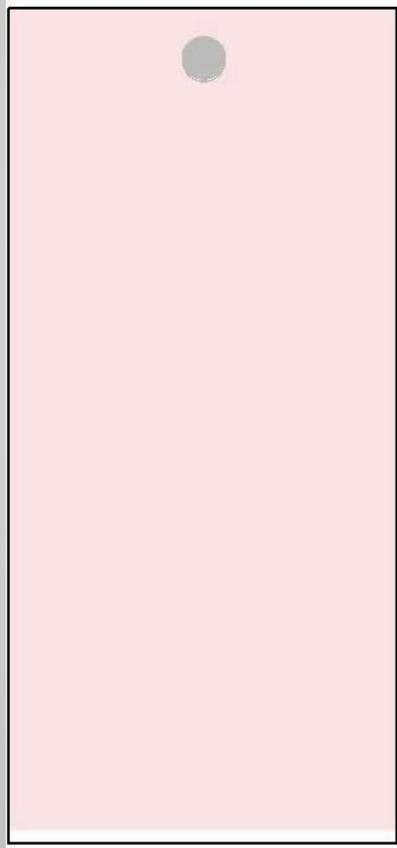
- Inertial waves
- Non-Boussinesq equations
- Non-Newtonian flows
- Particulate flows
- Gravity currents

Examples

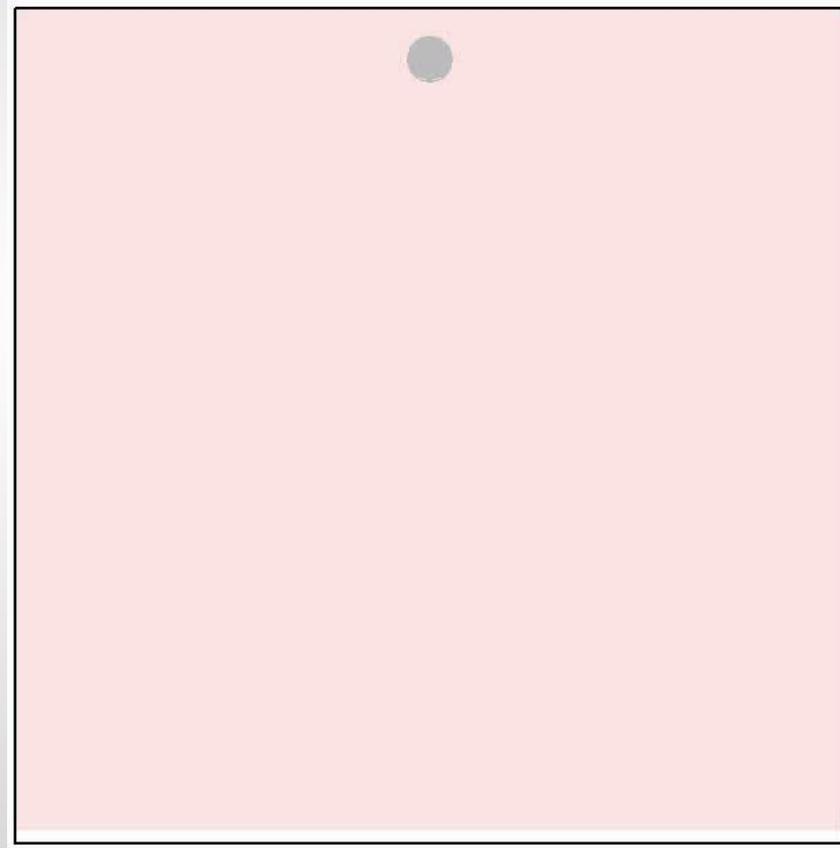
- Spin-up annulus
- Particulate flows
- Inertial waves

Particulate flows

No rotation*



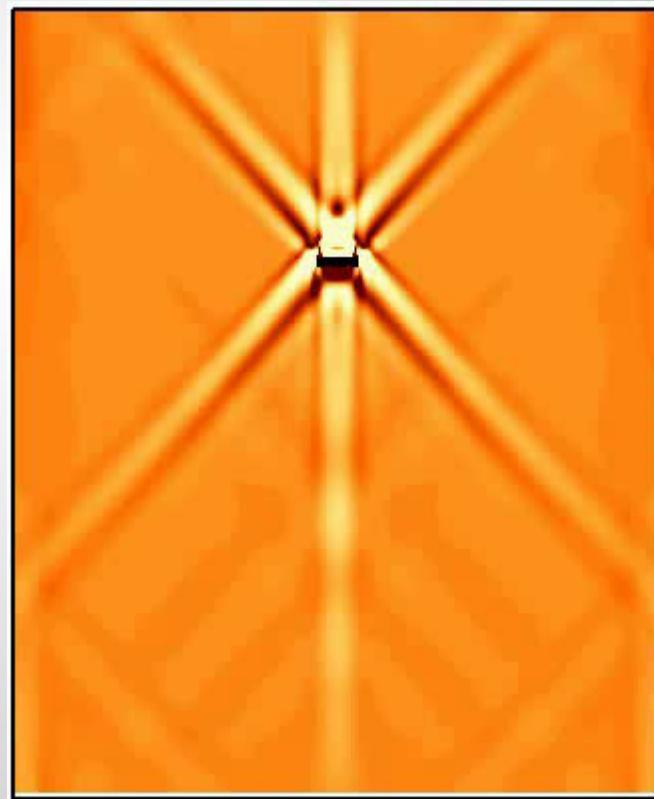
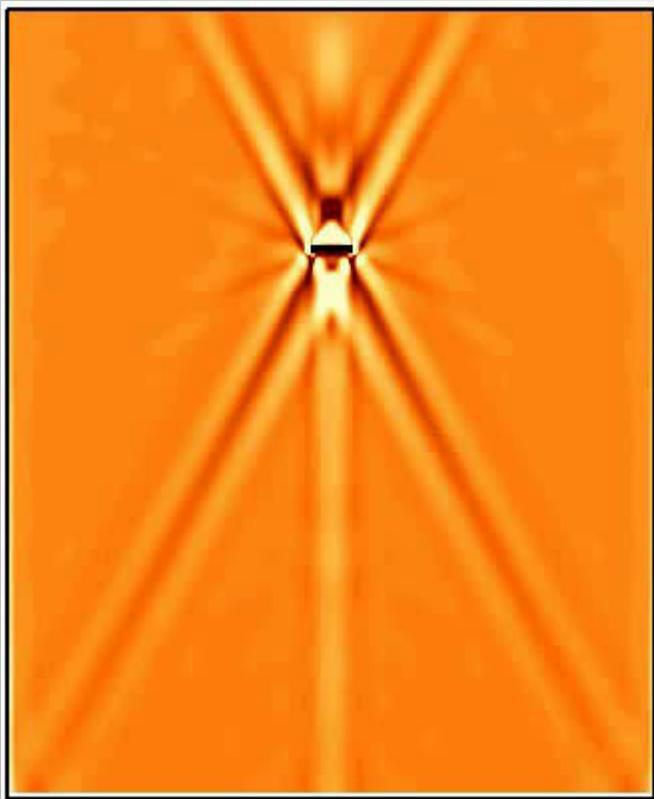
Rotation^



*Pacheco, J. R., Ruiz-Angulo, A., Zenit, R., and Verzicco, R., 'Fluid velocity fluctuations in a collision of a sphere with a wall,' *Physics of Fluids* **23**(6): 063301, 2011

^Pacheco, J. R. Lopez. J. M. Verzicco, R. 'Generation of inertial waves by a sphere settling in a rotating fluid.' *J. Fluid Mech.* 2012.

Inertial waves



Chaotic(?) mixing

